Aromaticity is one of the most important concepts in chemistry and refers to planar cyclic hydrocarbon molecules that exhibit delocalized $\pi$-bonds and unusual stability, such as benzene. This concept has been extended to metal-substituted organic molecules,[1,2] as well as main group organometallic complexes.[3–6] The recent discovery of aromaticity in all-metal clusters, for instance $[\text{Al}_4]^{2–}$ in $[\text{MAI}_4]^{–}$ (M = Cu, Li, Na),[7] has led to a flurry of research activities[8–16] and predictions of new aromatic metal clusters,[17–23] among which is a class of interesting cyclic species containing Cu.[20,23]

While main group clusters can give rise to $\sigma$- and $\pi$-aromaticity, transition-metal-containing clusters can exhibit d-orbital aromaticity or, more interestingly, $\delta$-aromaticity due to d–d bonding interactions. However, d-orbital aromaticity requires significant d–d bonding interactions. Unlike valent $s$ or $p$ orbitals, d orbitals are spatially more contracted, and their tendency to participate in chemical bonding depends strongly on the position of the transition metals in the periodic table and their coordination environments. The predicted aromaticity in the cyclic Cu-containing species,[20,23] for example, has not been verified experimentally. More promisingly, d-orbital aromaticity is expected to be found in early or 4d/5d transition metal systems, where strong d–d interactions are known. Here, we report experimental and theoretical evidence of d-orbital aromaticity in two early 4d and 5d transition metal oxide clusters, namely $[\text{M}_3\text{O}_9]^{–}$ and $[\text{M}_3\text{O}_9]^{2–}$ (M = W, Mo).
Early transition metal oxide clusters have recently gained experimental and theoretical attention due to their importance in catalysis.\cite{24–26} We investigated the structure and bonding of early transition metal oxide clusters using a combined photoelectron spectroscopy (PES) and theoretical approach with the aim of discovering specific structures or clusters that may mimic a catalytic site.\cite{27,28} Our experiment was carried out with a magnetic-bottle PES apparatus with a laser vaporization cluster source. A pure metal target (W or Mo) was used in the laser vaporization with a helium carrier gas containing 0.5% O₂. \[M_xO_y/C_0\]-type clusters with various stoichiometry were produced from reactions between the laser-induced plasma and O₂ in the carrier gas and were separated in a time-of-flight mass spectrometer. The \[M_3O_9/C_0\] clusters (M = W, Mo) of interest were selected and decelerated before being intercepted by a 157-nm laser beam from an F₂ excimer laser. Photoemitted electrons were collected by a magnetic bottle with nearly 100% efficiency and analyzed in a 3.5-m-long electron flight tube. The photoelectron spectra were calibrated against the known spectrum of Rh\[C_0\] and the apparatus had an electron energy resolution of \(\Delta E_k/E_k \approx 2.5\%\), i.e., approximately 25 meV for 1-eV electrons.

Figure 1 displays the 157-nm spectra of \[W_3O_9/C_0\] and \[Mo_3O_9/C_0\], which are similar, each with a weak and broad threshold band (X) followed by a large energy gap and more features at very high electron-binding energies. The \[M_3O_9\] clusters (M = W, Mo) of interest were selected and decelerated before being intercepted by a 157-nm laser beam from an F₂ excimer laser. Photoemitted electrons were collected by a magnetic bottle with nearly 100% efficiency and analyzed in a 3.5-m-long electron flight tube. The photoelectron spectra were calibrated against the known spectrum of Rh\[C_0\] and the apparatus had an electron energy resolution of \(\Delta E_k/E_k \approx 2.5\%\), i.e., approximately 25 meV for 1-eV electrons.

To understand the structures and bonding in the \[M_3O_9\] and \[M_3O_9/C_0\] clusters, we carried out density functional theory (DFT) calculations.\cite{29} Figure 2 displays the ground-state structures for \[M_3O_9\] and \[M_3O_9/C_0\]. The structures of neutral \[W_3O_9\] and \[Mo_3O_9\] are nearly identical, with \(D_3h\) symmetry, and each metal atom is bonded to two bridging O atoms and two terminal O atoms. The \[W_3O_9\] and \[Mo_3O_9\] anions also display \(D_3h\) symmetry, but the M–M distance is significantly reduced by 0.29 Å in \[W_3O_9\] and 0.25 Å in \[Mo_3O_9\] relative to the neutral structures, and this is accompanied by a large increase in the OMO angle and a large decrease in the MOM angle that makes the M₃O₉ ring almost a perfect hexagon in \[M_3O_9\] (Figures 2b and 2e). These geometrical changes are consistent with the broad ground-state PES band for the two clusters (Figure 1). Our calculated VDEs of 4.39 and 4.11 eV agree well with the measured values of 4.2 and 4.0 eV for \[W_3O_9\] and \[Mo_3O_9\], respectively, and our calculated adiabatic detachment energies (ADEs) for the two anions are 3.0 eV, which is smaller than the experimental threshold value of 3.5 eV. The large differences between the calculated ADEs and VDEs are consistent with the large difference in geometry between the anion and neutral ground-states. The smaller value of the calculated ADE relative to the experimental threshold value suggests that the 0–0 transition from the ground state of the anion to that of the neutral clusters may have negligible Franck–Condon factors due to the large difference in geometry, and the measured threshold detach-
ment energy can only serve as an upper limit for the true adiabatic value.

The surprisingly large geometry change induced by addition of a single electron to [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{2+} suggests that the molecular orbital that the extra electron occupies in the anion must involve strong metal–metal bonding. Figures 3a and 3b show the singly occupied molecular orbitals (SOMOs) for [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2+}, which indeed describe nearly pure three-center M–M o-bonding interactions with very little contribution from the O atoms. Judging by the huge M–M distance change induced by the single electron, these SOMOs must depict two strong three-center, one-electron (3c–1e) metal–metal bonds.

The SOMO of [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{-} is a nondegenerate a\textsubscript{1} orbital. Addition of a second electron to this a\textsubscript{1} orbital gives rise to closed-shell [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} dianions, which also display D\textsubscript{3h} symmetry (Figure 2c, f) and a three-center, two-electron (3c–2e) bond (Figures 3c and 3d). Our calculations show that the dianions are also unusually stable despite the strong intramolecular coulomb repulsion. In fact, [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} is predicted to be a stable gas-phase dianion with a calculated electron binding energy of 0.03 eV relative to [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{−}. The [Mo\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} dianion is predicted to be slightly unstable, with a small negative electron binding energy of −0.14 eV. More significantly, we found that the addition of a second electron to [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−} also induces a considerable shortening of the M–M bond length of 0.19 Å (0.20 Å) for [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} ([Mo\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−}) relative to [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−}.

To gain further insight into the 3c–1e and 3c–2e bonds in [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−} and [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3}, we estimated their resonance energies (see Supporting Information). Two isodesmic equations were used for [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} to eliminate the potential influence of the different electrostatic interactions in the doubly charged anions. When the intramolecular coulomb interactions of the dianions involved are balanced by the K\textsuperscript{+} cation a very large resonance energy of 1.05 eV is obtained, which is, in fact, comparable to the value estimated for benzene. This large resonance energy suggests that the [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} species may be viewed as highly aromatic. We further calculated their nucleus independent chemical shifts (NICS), which have been suggested as a magnetic criterion for aromaticity. We obtained NICS values of −21.5 and −20.5 at the ring center for [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} and [Mo\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3}, respectively (see Supporting Information), that are much higher than those predicted for the cyclic Cu-containing clusters[20,23] and are comparable to the multiply aromatic [Al\textsubscript{3}]\textsuperscript{3−}. The large resonance energies, the equal M–M bond lengths, and the large negative NICS values all suggest that [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} and [Mo\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} are highly aromatic species. More interestingly, we also obtained a sizable resonance energy for [W\textsubscript{3}O\textsubscript{9}]\textsuperscript{2−} (0.33 eV, see Supporting Information) even though it only contains a single delocalized electron, which suggests it may also be considered as aromatic; this was corroborated by the NICS values.

Trinuclear clusters of W and Mo have been studied extensively[32,33] However, all these cluster complexes involve M\textsuperscript{4+} metal centers, typified by the [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} core, whose metal–metal bonding can be described by the three classical 2c–2e metal–metal bonds: they are therefore not expected to exhibit aromaticity. Quasi-aromaticity has been invoked to explain the strong M–O interactions in the [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} and [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} clusters discussed here. Aromaticity has also been considered in trinuclear Pt\textsubscript{3} and Os\textsubscript{3} metal complexes solely on the basis of molecular orbital topologies[36,37] but it has not been verified using other criteria for aromaticity.

The [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−} and [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} clusters discussed here are unique in that they involve a single, fully delocalized metal–metal bond and may be considered as a new class of d-orbital aromatic molecules. The oxidation state of M can be considered to be fractional—+1/3 in [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−} and +16/3 in [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3}. Polyoxometalate clusters involving W and Mo are also well known.[38] The high stability of the [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{−} and [M\textsubscript{3}O\textsubscript{9}]\textsuperscript{16/3} clusters suggest that they may also be synthesized in the condensed phase. With the extra delocalized d-electrons, these clusters may exhibit novel chemical, electrochemical, and catalytic properties.

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Theoretical calculations were performed at the DFT level using the B3LYP hybrid functional (A. D. Becke, J. Chem. Phys. 1993, 98, 5648), augmented with two f-type and one g-type polarization functions for W (z(f) = 0.256, 0.825; z(g) = 0.627) for Mo and (z(f) = 0.338, 1.223; z(g) = 0.744), as recommended by Martin and Sundermann (J. M. L. Martin, A. Sundermann, J. Chem. Phys. 2001, 114, 3408), and the aug-cc-pVTZ basis set for oxygen (R. A. Kendall, T. H. Dunning, Jr., R. J. Harrison, J. Chem. Phys. 1992, 96, 6796). Scalar relativistic effects, i.e., the mass-velocity and Darwin effects, were taken into account by using the quasi-relativistic pseudo-potentials. Vertical electron detachment energies (VDEs) were calculated using the generalized Koopman’s theorem (D. J. Tozer, N. C. Handy, J. Chem. Phys. 1998, 109, 10180). The extra-fine integration grid was used to obtain accurate DFT results. The geometry optimizations and vibrational frequency analysis were accomplished with the NWChem 4.6 program (NWChem, A Computational Chemistry Package for Parallel Computers, Version 4.6: E. Apra et al., see Supporting Information) and the Molecular Science Computing Facility located at the Environmental Molecular Sciences Laboratory. Gaussian 3.0 was used to generate the three-dimensional contours of the calculated Kohn-Sham orbitals. Nucleus-independent chemical shifts (in ppm) were computed at the centers of the optimized structures using the GIAO-B3LYP (SDD for W and Mo, 6-31+G* for O) method with the Gaussian03 program (Gaussian03 (Revision B.04): M. J. Frisch et al., see Supporting Information).

Tozer, N. C. Handy, calculated using the generalized Koopman-HKkSs theorem (D. J. Tozer, N. C. Handy, J. Chem. Phys. 1998, 109, 10180). A number of possible structures were evaluated, among which those at 0.256, 0.825; ζ(g) = 0.627 for Mo and (ζ(f) = 0.338, 1.223; ζ(g) = 0.744), as recommended by Martin and Sundermann (J. M. L. Martin, A. Sundermann, J. Chem. Phys. 2001, 114, 3408), and the aug-cc-pVTZ basis set for oxygen (R. A. Kendall, T. H. Dunning, Jr., R. J. Harrison, J. Chem. Phys. 1992, 96, 6796). Scalar relativistic effects, i.e., the mass-velocity and Darwin effects, were taken into account by using the quasi-relativistic pseudo-potentials. Vertical electron detachment energies (VDEs) were calculated using the generalized Koopman’s theorem (D. J. Tozer, N. C. Handy, J. Chem. Phys. 1998, 109, 10180). The extra-fine integration grid was used to obtain accurate DFT results. The geometry optimizations and vibrational frequency analysis were accomplished with the NWChem 4.6 program (NWChem, A Computational Chemistry Package for Parallel Computers, Version 4.6: E. Apra et al., see Supporting Information) and the Molecular Science Computing Facility located at the Environmental Molecular Sciences Laboratory. Gaussian 3.0 was used to generate the three-dimensional contours of the calculated Kohn-Sham orbitals. Nucleus-independent chemical shifts (in ppm) were computed at the centers of the optimized structures using the GIAO-B3LYP (SDD for W and Mo, 6-31+G* for O) method with the Gaussian03 program (Gaussian03 (Revision B.04): M. J. Frisch et al., see Supporting Information).

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