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## “W-X-M” transformations in isomerization of $B_{39}^-$ borospherenes

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The Stone-Wales transformation plays an important role in the isomerization of fullerenes and graphenic systems. The continuous conversions between neighboring six- and seven-membered rings in the borospherene (all-boron fullerene)  $B_{40}$  had been discovered (Martínez-Guajardo *et al.* *Sci. Rep.* **5**, 11287 (2015)). In the first axially chiral borospherenes  $C_3 B_{39}^-$  and  $C_2 B_{39}^-$ , we identify three active boron atoms which are located at the center of three alternative sites involving five boron atoms denoted as “W”, “X”, and “M”, respectively. The concerted movements of these active boron atoms and their close neighbors between neighboring six- and seven-membered rings define the “W-X-M” transformation of borospherenes. Extensive first-principles molecular dynamics simulations and quadratic synchronous transit transition-state searches indicate that, via three transition states (TS1, TS2, and TS3) and two intermediate species (M1 and M2), the three-step “W-X-M” transformations convert the  $C_3 B_{39}^-$  global minimum into its  $C_2$  isomer at room temperature (300 K) and vice versa. The maximum barriers are only 3.89 kcal/mol from  $C_3$  to  $C_2 B_{39}^-$  and 2.1 kcal/mol from  $C_2$  to  $C_3 B_{39}^-$ , rendering dynamic fluxionalities to these borospherenes. Therefore, the “W-X-M” transformation plays an important role in the borospherenes and borospherene-based nanostructures. © 2016 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). [<http://dx.doi.org/10.1063/1.4954030>]

### I. INTRODUCTION

Shortly after the discovery of fullerenes,<sup>1</sup> the Stone-Wales transformation was proposed as a possible mechanism for interconversion between fullerene isomers.<sup>2</sup> By rotating a C-C bond by 90° with regard to the midpoint of the bond, the Stone-Wales transformation makes four hexagons turn into two pentagons and two heptagons. In particular, this transformation has the lowest formation energy in the isomerization of fullerenes and graphenic systems.<sup>3–9</sup>

The first all-boron fullerene  $D_{2d} B_{40}^{-/0}$ , referred to as borospherene in literature, was discovered in 2014 in a combined experimental and theoretical investigation,<sup>10</sup> and leads to a quick surge of borospherene chemistry, an area parallel to carbon fullerene chemistry.<sup>11–18</sup> The first axially chiral borospherenes,  $C_3$  and  $C_2 B_{39}^-$ , were observed in 2015.<sup>19</sup> Recently, three chiral  $B_{41}^+$ ,  $B_{42}^{2+}$ , and  $B_{44}$  cages were predicted using the extensive global-minimum searches and the first principle calculations.<sup>20,21</sup> All of these borospherenes are cubic-like in geometric structures which are composed of twelve interwoven boron double-chains (BDCs) with six  $B_6$  hexagons,  $B_7$  heptagons, and/or  $B_9$  nonagons on the faces.

Because of the electron deficiency of boron, there exist extensive multicenter B-B bonds in boron sheets and cages,<sup>10,19,22–26</sup> making their geometric structures fluxional. The structural fluxionalities of small boron clusters has been widely investigated to design nano-devices.<sup>18,27–32</sup> For instance, a series of Wankel motor molecules including  $B_{11}^-$ ,  $B_{13}^+$ ,  $B_{18}^{2-}$ , and  $B_{19}^-$  were discovered using the

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first-principles methods.<sup>28–32</sup> In particular, first-principles molecular dynamics (FPMD) simulations demonstrated that the structural transformation of B<sub>5</sub> fragment between six- and seven-membered rings is an element process in isomerization of the B<sub>40</sub> fullerene.<sup>18</sup>

The axially chiral B<sub>39</sub><sup>−</sup> fullerenes possess a distinctive structural fluxionality: B<sub>39</sub><sup>−</sup> hops between the C<sub>3</sub> and C<sub>2</sub> isomers with low energy barriers during FPMD simulations at 300K. In the two close-lying isomers, C<sub>2</sub> B<sub>39</sub><sup>−</sup> is the first borospherene observed with a tetracoordinate boron atom at the “defect” site of a boron double chain.<sup>19</sup> B<sub>39</sub><sup>−</sup> therefore may serve as a preferred model to investigate the structural transformations and isomerization of borospherenes.

We carefully investigated the structural characteristics of the B<sub>39</sub><sup>−</sup> fullerenes and found that both the C<sub>3</sub> and C<sub>2</sub> B<sub>39</sub><sup>−</sup> fullerenes can be viewed as the combinations of a spherical B<sub>36</sub> skeleton with three active boron atoms on the BDCs between neighboring hexagonal and heptagonal holes. The concerted movements of these active boron atoms and their directly bonded neighbors make the B<sub>39</sub><sup>−</sup> fullerene fluctuate between its low-lying isomers at room temperature.

## II. COMPUTATIONAL METHODS

The geometric structures of all B<sub>39</sub><sup>−</sup> fullerenes were optimized at the CAM-B3LYP/6-311+G\* and PBE0/6-311+G\* levels, respectively.<sup>33–38</sup> Both of their single point energies were calculated at the CCSD(T)/6-31G\* level<sup>39–42</sup> using the Molpro2012 package.<sup>43</sup> Though the optimized geometric structures of these borospherenes have only some negligible difference between the CAM-B3LYP and PBE0 functionals, the relative energies of CAM-B3LYP are more approximate to those of CCSD(T) than those of PBE0.<sup>19</sup> Therefore, the CAM-B3LYP functional was chosen to investigate the other properties of borospherenes. The transition states (TS) were searched using the quadratic synchronous transit (QST) approach<sup>44,45</sup> at the CAM-B3LYP/6-31G\* levels. To confirm the connectivity of minima and saddle points, the intrinsic reaction coordinates were scanned with the step size of 0.05 Bohr on both directions. All of these first principle calculations, except the CCSD(T) ones, were carried out using Gaussian 09 package.<sup>46</sup>

The FPMD simulations of the B<sub>39</sub><sup>−</sup> fullerenes were performed using the QUICKSTEP module of CP2K2.6 suite<sup>47</sup> at the temperatures of 200, 300, and 500K, respectively. The electronic part of calculations is handled with the BLYP functional<sup>48,36</sup> employing a mixed Gaussian and plane waves approach. The Gaussian basis set is double-zeta valence polarization basis set with Goedecker–Teter–Hutter pseudopotentials,<sup>49</sup> and the energy cutoff for plane waves is 280 Rydbergs. The simulations are launched from the optimized geometry of C<sub>2</sub> or C<sub>3</sub> B<sub>39</sub><sup>−</sup> fullerenes in a 1.5nm cubic box with random velocities assigned to the atoms, employing a Hoover thermal bath,<sup>50,51</sup> for a simulation time of 30 ps with 1.0 fs time steps. All B-B bond lengths above and under the active boron atoms were monitored along the trajectories to determine the sites of active boron atoms in the simulations.

## III. RESULTS AND DISCUSSION

It had been observed that the B<sub>39</sub><sup>−</sup> borospherenes have a C<sub>3</sub> cage global minimum with a close-lying C<sub>2</sub> cage isomer (Fig. 1(A) and 1(D)), both of which consist of a similar spherical B<sub>36</sub> skeleton and three active boron atoms on the BDCs at the waist. In the C<sub>3</sub> cage, all the three active boron atoms lie at the center of three upper sites mainly involving five boron atoms (denoted as “W”, Fig. 1(A)), while in the C<sub>2</sub> cage the three active boron atoms are located at the center of the lower (“M”), middle (“X”), and upper sites involving three B<sub>5</sub> units, respectively (Fig. 1(D)). Thus, each active boron atom in B<sub>39</sub><sup>−</sup> has three possible locations to occupy, at the center of either W, X, or M sites (Fig. 1).

Because the spherical B<sub>36</sub> skeleton, a new type of three-chain boron cage,<sup>52</sup> is relatively inactive, it can be used to construct a family of B<sub>39</sub><sup>−</sup> cages with three active boron atoms at its waist. Obviously, there are 3<sup>3</sup> = 27 combinations of the “W”, “X”, and “M” active boron atoms in this B<sub>39</sub><sup>−</sup> family, whose specific geometric structures can be systematically denoted by three occupied sites of

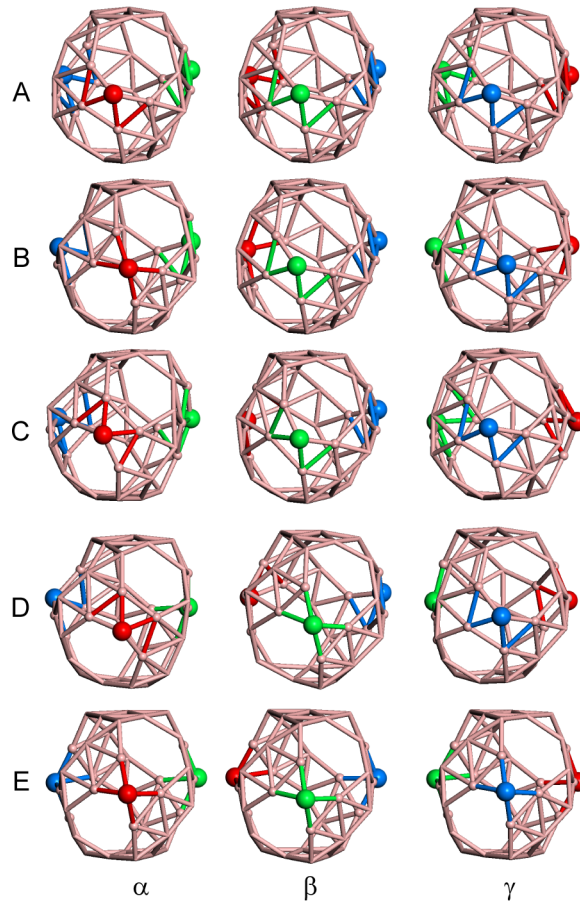


FIG. 1. Typical geometric structures of the  $B_{39}^-$  fullerenes: (A)  $C_3$ , (B)  $M_1$ , (C)  $M_2$ , (D)  $C_2$ , and (E)  $D_3$ . The  $\alpha$ ,  $\beta$ , and  $\gamma$  are their front, right-back, and left-back views, and the red, green, and blue balls indicate three active boron atoms at the waist, respectively.

active boron atoms (TABLE I). For instance, the  $C_3$  and  $C_2$   $B_{39}^-$  fullerenes have the “WWW” and “MXW” structures (Fig. 1(A) and 1(D)), respectively.

In this family, all 27  $B_{39}^-$  cages are classified into six nonequivalent groups:  $D_3$ ,  $C_3$ ,  $C_2$ ,  $M_1$ ,  $M_2$ , and  $M_3$ , which include 1, 2, 6, 6, 6, and 6 specific geometric structures, respectively (TABLE I). The high symmetry  $D_3$  “XXX”  $B_{39}^-$  fullerene should be taken as the primitive cage of the  $B_{39}^-$  family: its distortions will produce all other isomers with  $C_3$ ,  $C_2$ , or  $C_1$  symmetries. The  $M_3$  isomers

TABLE I. Specific geometric structures and relative energies (kcal/mol) of the  $B_{39}^-$  fullerenes.

Name	Symmetry	Geometric structures	$E_{CAM-B3LYP}$	$E_{PBE0}$	$E_{CCSD(T)//CAM-B3LYP}$	$E_{CCSD(T)//PBE0}$
$D_3$	$D_3$	XXX	4.62	7.04	4.64	10.31
$C_3$	$C_3$	WWW	0.00	0.00	0.00	0.00
		MMM				
$M_1$	$C_1$	XWW, WXW, WWX	2.84	4.57	3.78	3.97
		XMM, MXM, MMX				
$M_2$	$C_1$	MWW, WMW, WWM	2.68	4.67	3.98	3.90
		WMM, MWM, MMW				
$C_2$	$C_2$	MXW, WMX, XWM	0.73	3.07	1.74	2.20
		WXM, MWX, XMW				
$M_3$	-	XXW, XWX, WXX	-	-	-	-
		XXM, XMW, MXX				

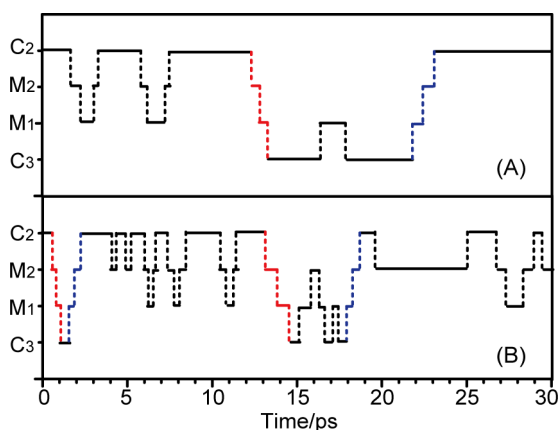


FIG. 2. Presences (horizontal solid lines) and transformations (vertical dashed lines) of  $C_3$ ,  $M_1$ ,  $M_2$ , and  $C_2$   $B_{39}^-$  fullerenes at 300K (A) and 500K (B).

are not true local minima. The energies of the  $D_3$ ,  $M_1$ , and  $M_2$  isomers are only 4.64, 3.78, and 3.98 kcal/mol less stable than the global minimum  $C_3$  cage at the CCSD(T)/6-31G\* level. These newly predicted species are expected to coexist with the observed  $C_3$  and  $C_2$  isomers in experiments under certain conditions (TABLE I), although they were not identified in the previously reported photoelectron spectroscopy (PES) of  $B_{39}^-$  which is rather congested and may contain contributions from other isomers.<sup>19</sup>

To locate all possible coexisting borospherenes, extensive FPMD simulations of the  $C_2$  and  $C_3$   $B_{39}^-$  fullerenes were performed at 200, 300, and 500K, respectively.<sup>19,53</sup> At 200K, both of the  $C_2$  and  $C_3$   $B_{39}^-$  fullerenes are stable and cannot be converted into the other low-lying isomers. However, at 300K and especially 500K, the  $C_3$ ,  $C_2$ ,  $M_1$ , and  $M_2$   $B_{39}^-$  isomers all coexist in the FPMD simulations of  $B_{39}^-$ , except the  $D_3$  one with an XXX sequence (Fig. 2).

In the FPMD simulations at 300K and 500K, there exist three-step reversible structural transformations among the  $C_3$ ,  $C_2$ ,  $M_1$ , and  $M_2$   $B_{39}^-$  isomers (FIG. 2): from  $C_3$  to  $M_1$  (FIG. 1(A $\alpha$ ) and 1(B $\alpha$ )), from  $M_1$  to  $M_2$  (Fig. 1(B $\alpha$ ) and 1(C $\alpha$ )), and from  $M_2$  to  $C_2$  (Fig. 1(C $\beta$ ) and 1(D $\beta$ )). All of these transformations originate from the movements of three active boron atoms between different sites. For an active boron atom, its “W-X” transformation expands the lower hexagon into a heptagon by breaking a B-B bond under it, and its “X-M” transformation shrinks the upper heptagon into a hexagon by forming another B-B bond above it (Fig. 3), which was firstly discovered in isomerization of the  $B_{40}$  fullerene.<sup>18</sup> Obviously, an active boron atom has only two elementary transformations: “W-X” and “X-M”. Here, the “W-X-M” transformation is used to denote all elementary transformations and their various combinations of active boron atoms. Such a mechanism can be easily extended to other borospherenes involving octagonal and nanogon rings etc, depending on the specific geometries of the systems.<sup>20,21</sup>

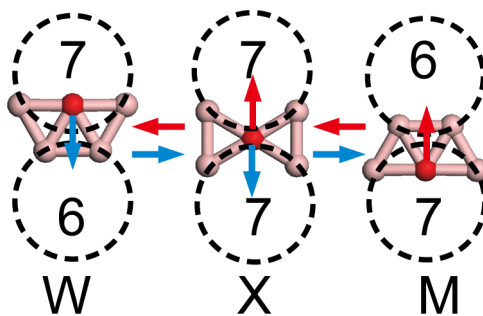


FIG. 3. “W-X-M” transformation between neighboring hexagon (6) and heptagon (7). The blue and red arrows denote the forward and backward “W-X-M” transformation, respectively.



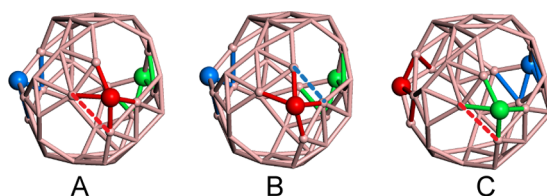


FIG. 4. Typical geometric structures of three transition states in the isomerization between the  $C_3$  and  $C_2$   $B_{39}^-$  fullerenes: (A)  $TS_1$  (between  $C_3$  and  $M_1$ ), (B)  $TS_2$  (between  $M_1$  and  $M_2$ ), and (C)  $TS_3$  (between  $M_2$  and  $C_2$ ). The red, blue, and green balls indicate three active boron atoms, and the red and blue dashed lines indicate the broken and forming B-B bonds in the reactions, respectively.

In the FPMD simulations at 300 and 500K, the isomerization from the global minimum  $C_3$  cage to its close-lying  $C_2$  isomers (blue and red lines in Fig. 2) includes the “W-X” and “X-M” transformations of an active boron atom, and the “W-X” transformation of another active boron atom. Furthermore, the isomerization between the  $C_3$  “WWW” and “MMM” structures can also occur by three successive “W-X-M” transformations of three active boron atoms:  $C_3(\text{WWW}) \rightarrow M_1(\text{XWW}) \rightarrow M_2(\text{MWW}) \rightarrow C_2(\text{MXW}) \rightarrow M_2(\text{MMW}) \rightarrow M_1(\text{MMX}) \rightarrow C_3(\text{MMM})$ . In addition, there are also many incomplete isomerization processes, such as the  $C_2$ - $M_2$ - $C_2$ ,  $C_3$ - $M_1$ - $C_3$ , and  $C_3$ - $M_1$ - $M_2$ - $M_1$ - $C_3$  transformations.

Beside the  $B_{39}^-$  borospherenes, there is the similar “W-X-M” transformation between the six- and seven-membered rings of  $B_{40}$  borospherene in the FPMD simulations at high temperature (1200K).<sup>13</sup> Therefore, the “W-X-M” transformation plays an important role in the isomerization of borospherenes, just as the role of Stones-Wales transformation in the fullerenes and graphenic systems. However, in contrast with the C-C bond rotation in Stones-Wales transformation, the “W-X-M” transformation originates from the movement of some active boron atoms. In addition, because of the axial chirality of these  $B_{39}^-$  fullerenes, there is another family of degenerate enantiomers, in which the “W-X-M” transformations can also occur in their isomerization.

The most important isomerization is the reaction between the  $C_3$  and  $C_2$   $B_{39}^-$  fullerenes, whose transition states were found between  $C_3$  and  $M_1$  ( $TS_1$ ), between  $M_1$  and  $M_2$  ( $TS_2$ ), and between  $M_2$  and  $C_2$  ( $TS_3$ ) (Fig. 4 and 5). The scans of their intrinsic reaction coordinates confirmed their reaction mechanisms.<sup>54</sup> Their small imaginary harmonic vibrational frequencies ( $TS_1$ :  $104i$ ,  $TS_2$ :  $55i$ , and  $TS_3$ :  $89i$   $\text{cm}^{-1}$ ) demonstrated that the structural transformations are moderate.

The energies of these transition states are 3.89, 3.60, and 2.84 kcal/mol with respect to the  $C_3$  global minimum at the CAM-B3LYP/6-311+G\* level, respectively. Then their barriers from  $C_3$  to  $C_2$  isomers are 3.89, 0.76, and 0.16 kcal/mol so that the  $C_3$  to  $M_1$  transformation is the rate determining step. The maximum barrier of 3.89 kcal/mol implies that the isomerization from  $C_3$  to  $C_2$  isomers could occur at the room temperature, just as shown in the FPMD simulation at 300K. Reversibly, the barriers from  $C_2$  to  $C_3$  isomers are 2.11, 0.92, and 1.05 kcal/mol at the

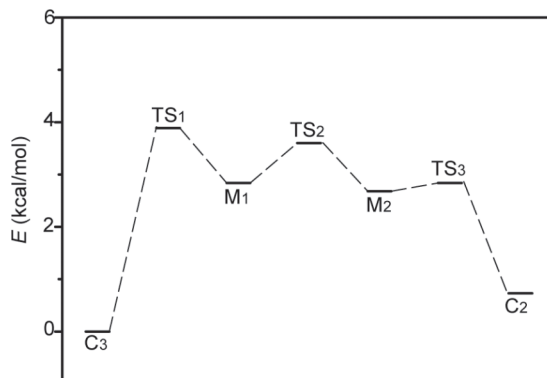


FIG. 5. Reaction pathway between the  $C_3$  and  $C_2$   $B_{39}^-$  fullerenes.

CAM-B3LYP/6-311+G\* level, respectively, which is much lower than the maximum barrier (3.89 kcal/mol). Therefore, the isomerization from  $C_2$  to  $C_3$  isomers is easier than that from  $C_3$  to  $C_2$  isomers.

The maximum isomerization barrier (3.89 kcal/mol) of  $B_{39}^-$  fullerenes is much lower than that of the Stone-Wales transformation in carbon materials (162.9 kcal/mol at the B3LYP/6-311G\* level).<sup>3</sup> Obviously, the isomerization of  $B_{39}^-$  boron fullerenes are easier than that of carbon fullerenes. Furthermore, this barrier is also lower than that of isomerization (14.3 kcal/mol) in the  $B_{40}$  fullerene.<sup>18</sup> Compared with the one-step isomerization of the  $B_{40}$  fullerene, there are three-step successive “W-X-M” transformations between two coexisting  $C_3$  and  $C_2$   $B_{39}^-$  isomers. Therefore many  $B_{39}^-$  cage isomers coexist at room temperature, which provides a reaction kinetic origin of the polymorphism of boron fullerenes.<sup>53</sup>

#### IV. CONCLUSIONS

Based on the structural characteristic of  $C_3$  and  $C_2$   $B_{39}^-$  borospherenes, a family of  $B_{39}^-$  fullerenes were discovered, all of which consist of a spherical  $B_{36}$  skeleton and three active boron atoms. Each of the active boron atoms has three alternative sites (denoted as “W”, “X”, and “M”). All of the stable  $B_{39}^-$  fullerenes can coexist and transform each other at 300 and 500 K by the “W-X-M” transformations, which was firstly discovered in isomerization of the  $B_{40}$  fullerene. These “W-X-M” transformations play an important role in the isomerization of borospherenes, just as the Stone-Wales transformation in isomerization of carbon materials. In contrast with the C-C bond rotation in Stone-Wales transformation, the “W-X-M” transformation originates from the movement of some active boron atoms. The maximum barrier is only 3.89 kcal/mol from the  $C_3$  to  $C_2$   $B_{39}^-$  fullerenes so that they can transform each other easily at room temperature. This finding reveals the dynamic characteristic of borospherenes and can be used to design the new borospherene-based materials.

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